implies that

$$
W_{k i}^{(0)}=A_{i} \delta_{k i},
$$

where usually we think of $A_{i}$ depending only on $x$. Here $A_{i}$ also depends on both $\zeta$ and $t$ (for example, the time dependence determined via (1.2) could imply $A_{2}$ $=\exp \left(i \zeta^{3} t\right)$ ). By considering the next order equation, $O(1)$, the $x$ dependence of $A_{i}$ may be obtained:

$$
\begin{aligned}
& i\left(d_{k}-d_{i}\right) W_{k i}^{(1)} \\
& \quad=\frac{\partial}{\partial x} W_{k i}^{(0)}-\sum_{m} N_{k m} W_{m i}^{(0)}=\frac{\partial A_{1}}{\partial x} \delta_{k i}-A_{i} N_{k t} .
\end{aligned}
$$

Since $N_{i i}=0, \partial A_{i} / \partial x=0$ and

$$
W_{k i}^{(1)}= \begin{cases}0, & k=i, \\ -N_{k i} A_{i} / i\left(d_{k}-d_{i}\right), & k \neq i,\end{cases}
$$

where $W_{i i}^{(1)}=0$, because $W_{i i}^{(1)} \neq 0$ is equivalent to the introduction in $A_{1}$ of $O(1 / \zeta)$ term. In a similar way, the higher order terms may be directly calculated. Note that the equations for $W_{k i}^{(p)}$ will be linear. From these results we see that

$$
V_{k i}=A_{i}(\zeta, t) \exp \left(i \zeta d_{t} x\right)\left[\delta_{k i}+O(1 / \zeta)\right] .
$$

By considering the variables $\Gamma_{k t}=V_{k i} / V_{i t}$, the as yet unknown expression $A_{i}(\zeta, t)$ will be of no importance.
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# Addendum: A classical perturbation theory [J. Math. Phys. 18, 110 (1977) 

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