

Physics 7B Sections 103/107 (Bordel)

Date: March 14-16, 2022

Administrative Notes:

Wednesday is discussion again, no actual lab.
Held remotely again at this zoom link.

Lab grades posted.

MT2 covers material from electric point charges
up to and including DC circuits.

Potential

$V(x, y, z)$ - Scalar (SI Volt)

Suppose you have $V(x, y, z)$.

1) Electric Field

$$V(x, y, z) \rightarrow \vec{E}(x, y, z)$$

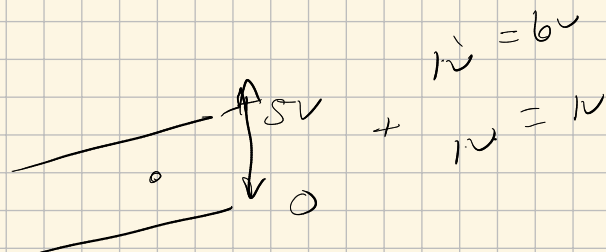
$$\vec{E} = -\vec{\nabla} V(x, y, z)$$

$$1D): \vec{E}(x) = -\frac{dV}{dx} \hat{i}$$

$$3D): \vec{E}(x, y, z) = -\frac{\partial V}{\partial x} \hat{i} - \frac{\partial V}{\partial y} \hat{j} - \frac{\partial V}{\partial z} \hat{k}$$

$$V(x, y, z) \rightarrow \vec{E}(x, y, z)$$

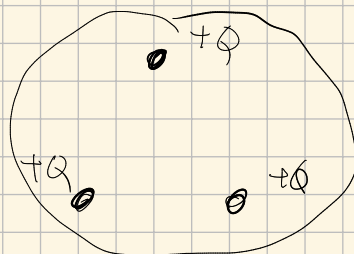
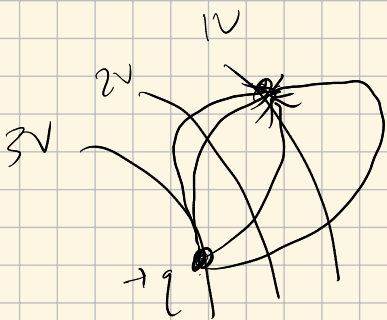
$$V(x, y, z) + V_0 \rightarrow$$



$$2) W = q(\Delta V) = q(v_f - v_i)$$

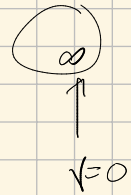
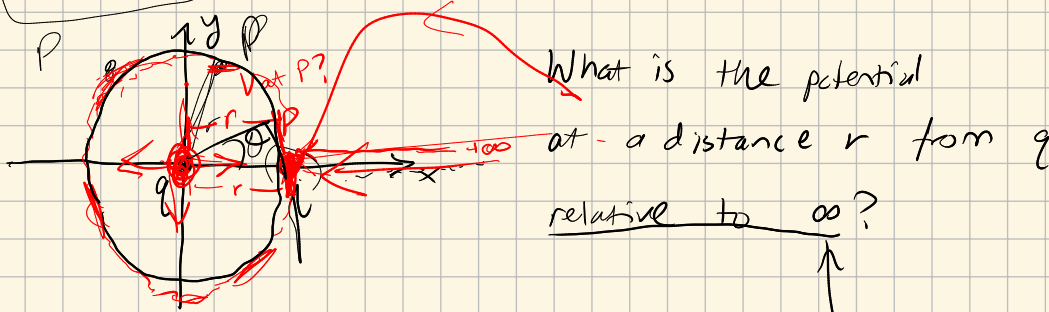
$$\frac{\Delta V}{\Delta x}$$

path-independent



Electric field \rightarrow potential

$\vec{E}(x,y,z) \rightarrow V(x,y,z)$



at ∞ , $V=0$

$$\Delta V = - \int_{+\infty}^r \vec{E} \cdot d\vec{l} = - \int_{+\infty}^r k \frac{q}{x^2} \hat{i} \cdot d\vec{x}$$

$$= - \left[-k \frac{q}{x} \right]_{+\infty}^r = k \frac{q}{r} - \cancel{k \frac{q}{\infty}} = k \frac{q}{r} \quad \checkmark$$

$V(r)$

$$V(r) - V(\infty) = k \frac{q}{r}$$

\parallel

$$V(r) = k \frac{q}{r}$$

for a point charge

$$\vec{E} \rightarrow V$$

Problem #1

Potential of a point charge is $k \frac{q}{r}$

Suppose you have a circle of radius R carrying uniform linear charge density λ . What is the electric field at a height z above the center?



Approach #1:

$$\vec{E} = \int d\vec{E} = \int d\vec{E}_x + \int d\vec{E}_y + \int d\vec{E}_z$$

$|\vec{E}|$ (trig function)

Approach #2:

Potential also obeys superposition.

$$V_{\text{total}}(z) = \int dV$$

$$= \int k \frac{dq}{\sqrt{R^2 + z^2}} = \frac{k}{\sqrt{R^2 + z^2}} \int dq$$

$Q_{\text{total}} = \lambda 2\pi R$

$$= \frac{k 2\pi R \lambda}{\sqrt{R^2 + z^2}}$$

$$V(z) \rightarrow \vec{E}(z)?$$

$$\vec{E}(z) = -\vec{\nabla} V(z) = -\frac{\partial V}{\partial x} \hat{i} - \frac{\partial V}{\partial y} \hat{j} - \frac{\partial V}{\partial z} \hat{k}$$

$$\vec{E}(z) = - \frac{\partial V}{\partial z} \hat{k} = \left(\frac{dV}{dz} \right) \hat{k}$$

$$V(z) = k \frac{2\pi R \lambda}{\sqrt{R^2 + z^2}}$$

$$\frac{dV}{dz} = k \frac{2\pi R \lambda}{(R^2 + z^2)^{3/2}} \cdot \left(-\frac{1}{2}\right) (2z) = -k \frac{2\pi R \lambda z}{(R^2 + z^2)^{3/2}}$$

$$\vec{E}(z) = k \frac{2\pi R \lambda z}{(R^2 + z^2)^{3/2}} \hat{k}$$

When there is symmetry, and enough of it, you can probably use Gauss's law to find the field

But otherwise, you might want to find the potential first and then compute its gradient to find the field.

Potential ° SI
 Volts = $\frac{N \cdot m}{C} = \frac{J}{C}$

Electric field ° °

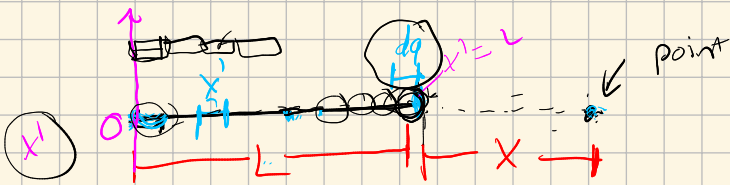
$$N/C = V/m$$

($F = qE$)

($E = -\nabla V$...)

Problem

$$\sum_{i=0}^{\infty} \frac{q_i}{r_i^2}$$



linear
✓

A line of length L carries uniform charge density λ . What is the potential at a distance x from the end as shown above?

Assume $x > 0$.

$V(x)?$

$$1) \quad V(x) = \int dV$$

$$= \int k \frac{dq}{x+L-x'}$$

$$= \int_0^L k \frac{\lambda dx'}{x+L-x'}$$

$$= k\lambda \left[-\ln(x+L-x') \right]_0^L$$

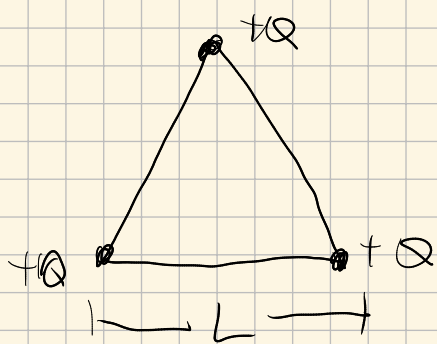
$$= k\lambda \left[-\ln(x) + \ln(x+L) \right]$$

$$= k\lambda \ln\left(\frac{x+L}{x}\right) \quad \checkmark$$

$$\int_0^L \lambda dx' = \lambda L = Q$$

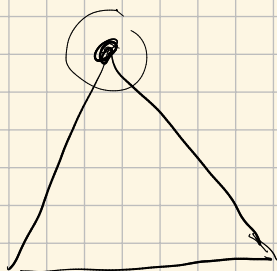
Problem

Consider 3 point charges all of value $+Q$. What is the work required to bring them from infinitely far apart to sit on the vertices of an equilateral triangle of side length L ?



1st charge :

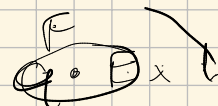
$$W_1 = ?$$



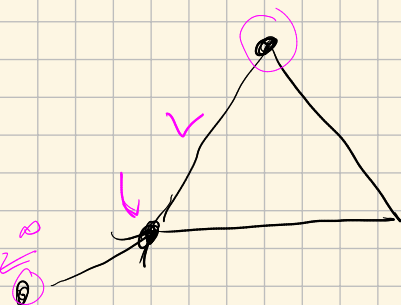
$$W = Q \Delta V = 0$$

2nd charge

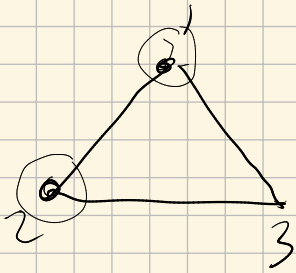
created by first charge



$$W_2 = Q \Delta V = Q \left(\underbrace{V_f}_{k \frac{Q}{L}} - \underbrace{V_i}_0 \right) = k \frac{Q^2}{L}$$



3rd charge



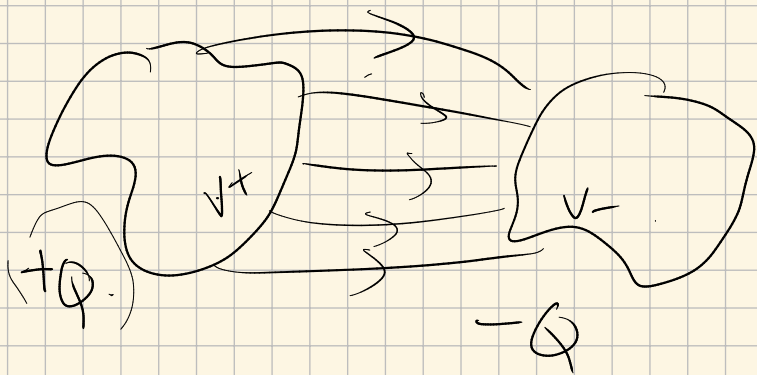
$$\left. \begin{aligned} V_{13} &= k \frac{Q}{L} \\ V_{23} &= k \frac{Q}{L} \end{aligned} \right\} V_f = k \frac{2Q}{L}$$

$$V_i = 0$$

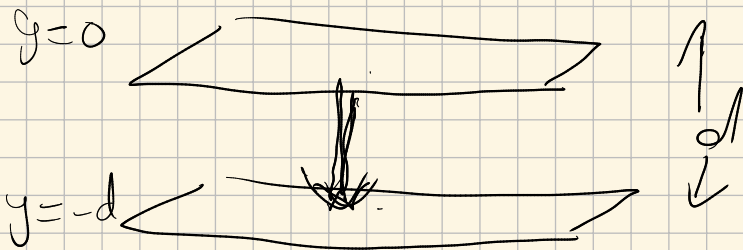
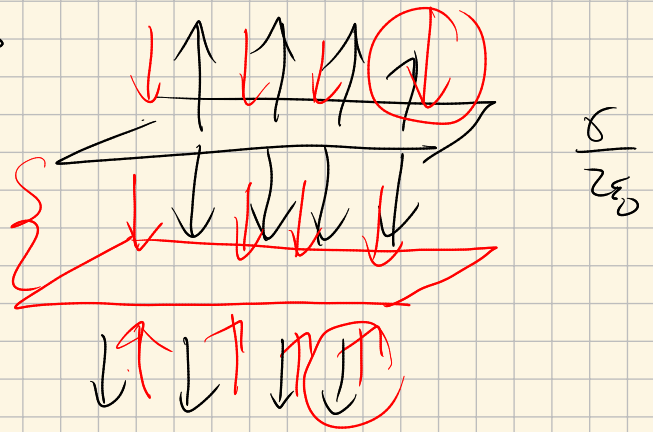
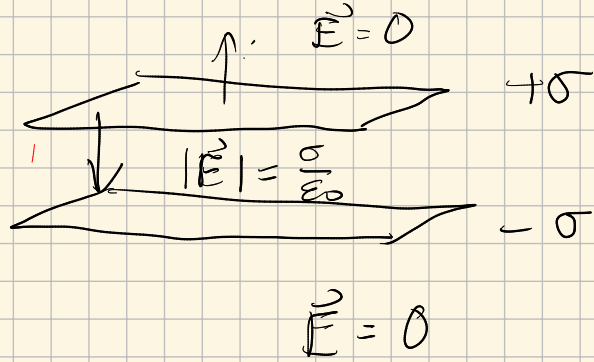
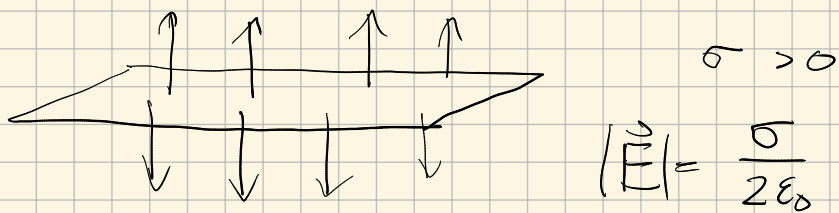
$$W_3 = Q (V_f - V_i) = k \frac{2Q^2}{L}$$

$$W_{\text{total}} = W_1 + W_2 + W_3 = 0 + k \frac{Q^2}{L} + k \frac{2Q^2}{L} = \boxed{3k \frac{Q^2}{L}}$$

Capacitors



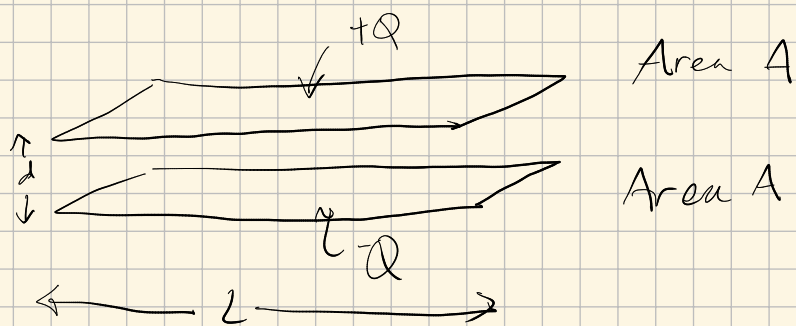
$$\frac{Q}{\Delta V} = C \quad (\text{F})$$



$$\vec{E} = -\frac{\sigma}{\epsilon_0} \hat{j}$$

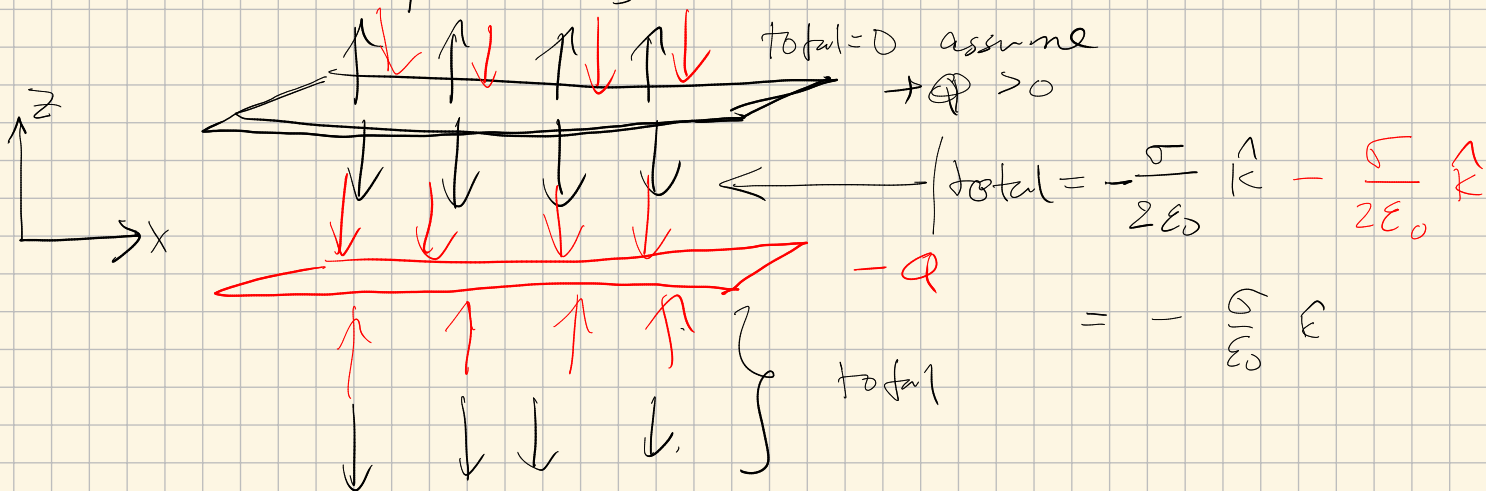
$$\Delta V = -\int \vec{E} \cdot d\vec{l} = -\int_0^{-d} \left(-\frac{\sigma}{\epsilon_0} \hat{j}\right) \cdot (-d\hat{j})$$

$$= \frac{\sigma}{\epsilon_0} d$$



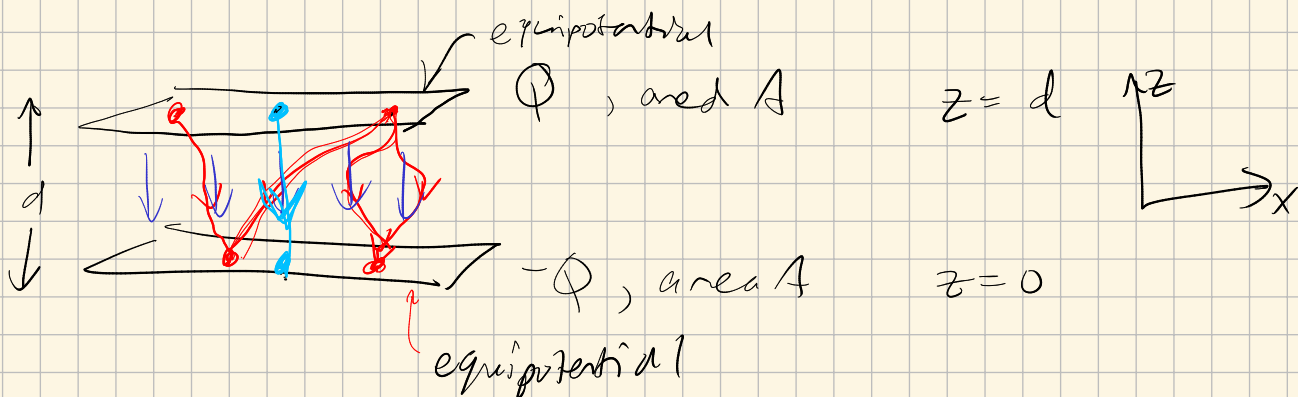
1) what does \vec{E} look like?

when $d \ll L$, approximate this as 2 infinite sheets



electric field of ∞ sheet has constant $|\vec{E}|$ everywhere

\vec{E} field in between two oppositely charged infinite sheets will be $-\frac{\sigma}{\epsilon_0} \hat{k}$



- 1) Imagine putting charge $+Q$ on one and $-Q$ on the other
- 2) Determine \vec{E} between them
- 3) ΔV between the two conductors.

$$\vec{E} = -\vec{\nabla} V \rightarrow \Delta V = - \int_a^b \vec{E} \cdot d\vec{l}$$

finite plate \rightarrow infinite one
charge

Approx: finite plate of area A .

assume the charge is uniform ;

if \rightarrow put $+Q$ on one plate and $-Q$ on the other \rightarrow infinite sheet.

$$E = \frac{\sigma}{\epsilon_0} = \frac{Q}{A \epsilon_0}$$

$$\Delta V = - \int_d^0 \vec{E} \cdot d\vec{l} = - \frac{Q}{A \epsilon_0} \int_d^0 dz = - \frac{Q}{A \epsilon_0} (0-d)$$

$$= \frac{Q}{A \epsilon_0} d$$

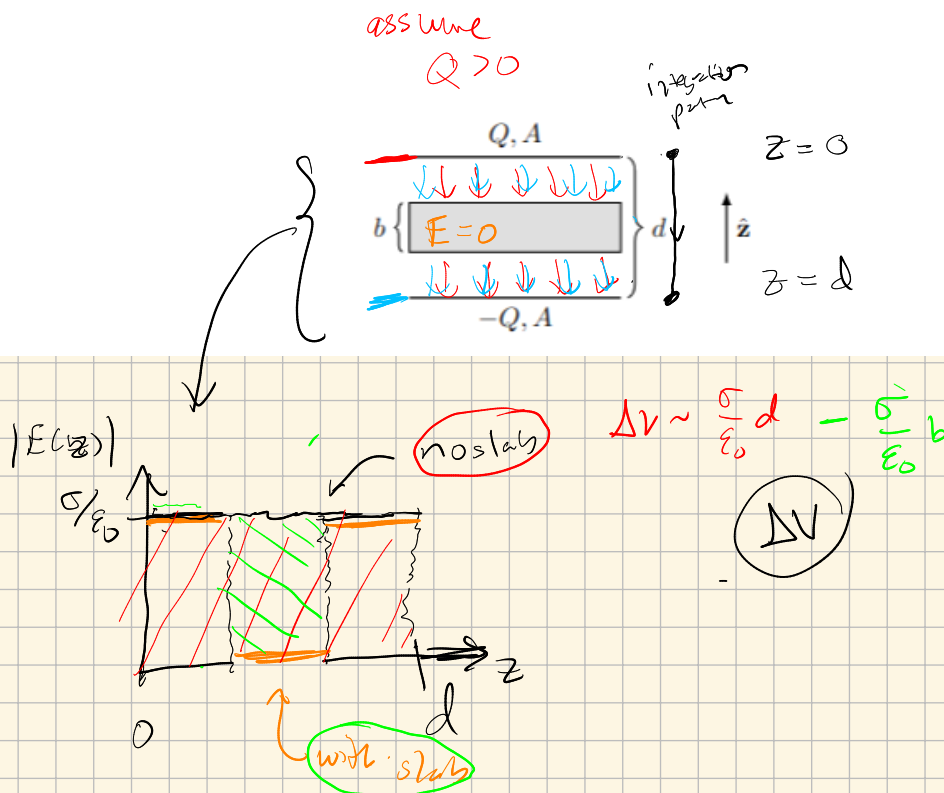
$$C = \frac{Q}{\Delta V} = \frac{Q}{\frac{Q}{A \epsilon_0} d} = \frac{A}{d} \epsilon_0 \quad (F)$$

• Problem 3: A conductor inside a capacitor [3]

This is problem 7 of WS E6 in your course workbook.

Consider a parallel-plate capacitor with plates of area A separated by some small distance d . It is initially charged up to some charge Q , and then disconnected from the battery. A neutral conducting metal slab, also of area A and with width $b < d$, is then inserted between the plates as shown below. What is the capacitance of this configuration? (Hint: it may help to start by thinking about what the electric field looks like inside the capacitor.)

¹We haven't specifically talked about conductors in any worksheet yet, but it's worth noting that conductors have some special properties. Most importantly, $\mathbf{E} = 0$ in the bulk of any conductor; any net charge on a conductor resides on the surface; and the surface of a conductor is always an equipotential. Review Section 21.9 in your textbook for more details [2].



- 1) Imagine charge $+Q$ on one conductor, $-Q$ on the other
- 2) \vec{E} -field
- 3) ΔV between the 2 conductors
- 4) $C = Q/\Delta V$

$$\Delta V = - \int \vec{E} \cdot d\vec{\ell} = - \int_0^d |E(z)| dz = - \left(\frac{\sigma}{\epsilon_0} (d-b) \right)$$

area under the curve $|E(z)|$

$$\sigma = \frac{Q}{A}$$

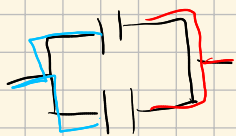
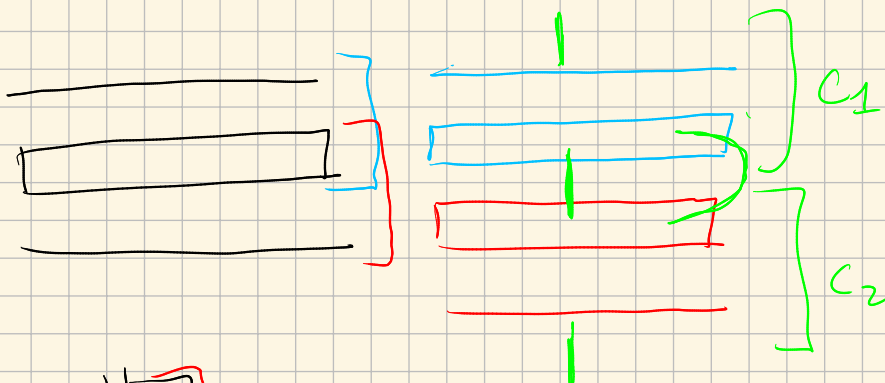
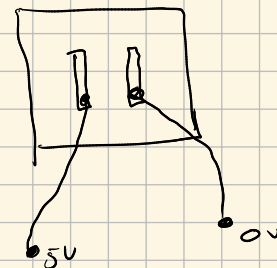
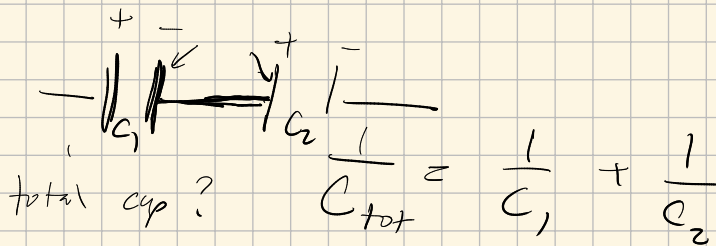
$$- \left(\frac{Q}{A \epsilon_0} (d-b) \right)$$

$$\Delta V = - \left(\frac{Q}{A \epsilon_0} (d-b) \right)$$

$$b < d$$

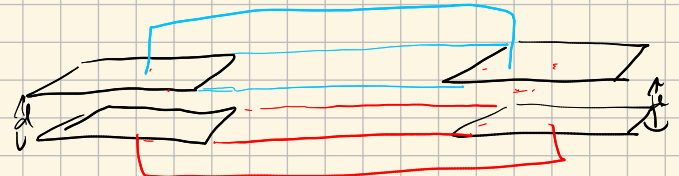
$$C = \frac{Q}{|\Delta V|} = \frac{Q}{\left| - \left(\frac{Q}{A \epsilon_0} (d-b) \right) \right|} = \frac{A \epsilon_0}{d-b}$$

Capacitors in series

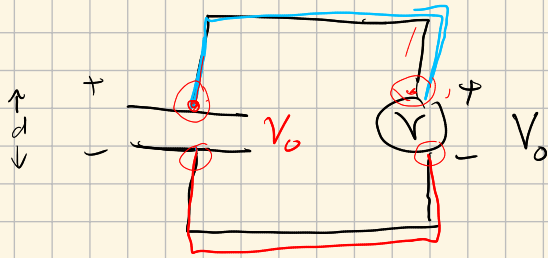


$$C_{tot} = C_1 + C_2$$

$$C = \frac{2A}{d} \epsilon_0$$



Suppose the capacitor has initial separation d , and is connected to voltage source (e.g. battery) and we pull the plates apart to distance $2d$.



$$Q = CV$$

What changes and by how much?

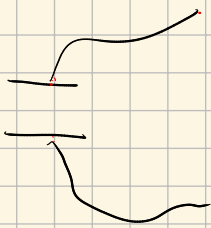
$$C = \epsilon_0 \frac{A}{d}$$

$$Q \rightarrow \frac{1}{2} Q$$

$$C \rightarrow \frac{1}{2} C$$

$$V \rightarrow V$$

Suppose we disconnect the battery before separating the plates to $2d$:



$$C = \epsilon_0 \frac{A}{d}$$

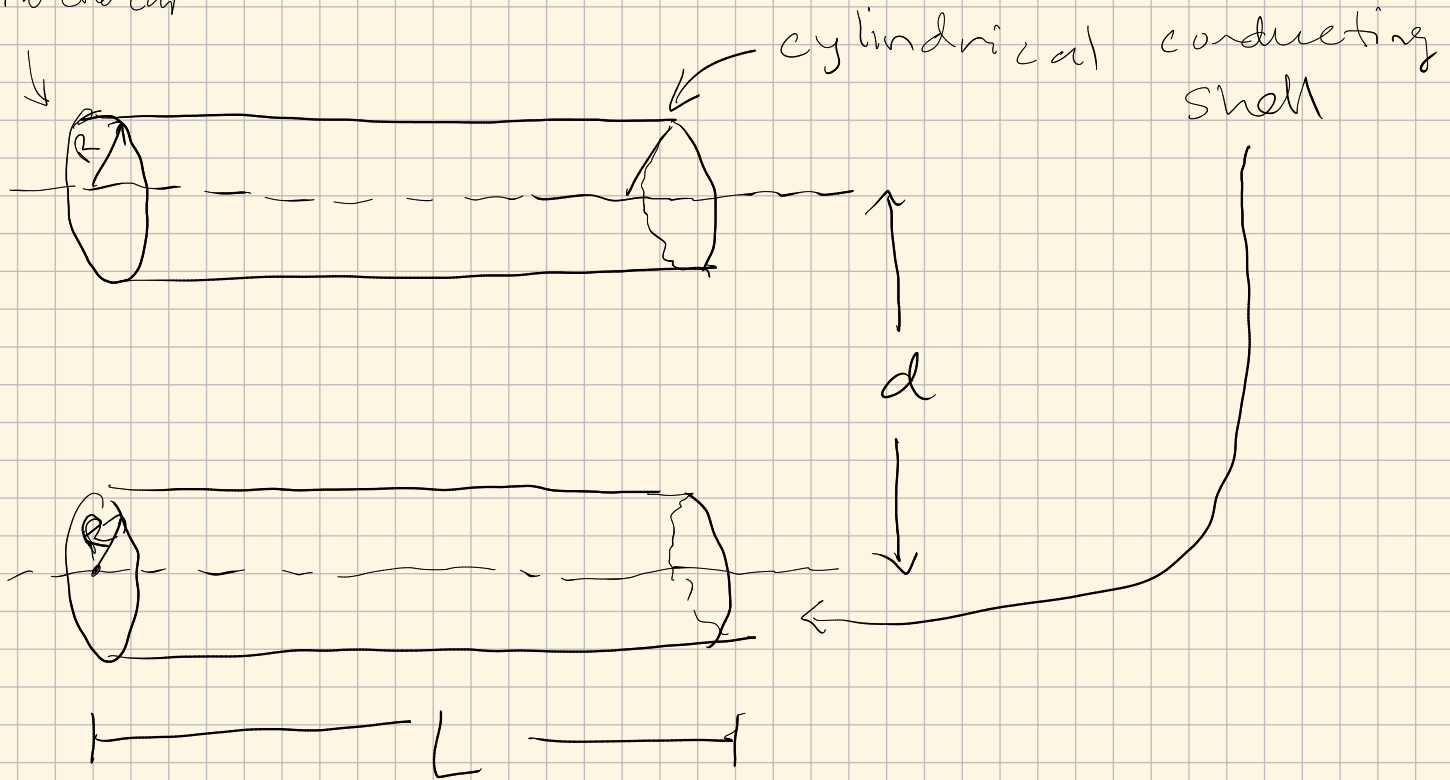
$$Q \rightarrow Q$$

$$C \rightarrow \frac{1}{2} C$$

$$V \rightarrow 2V$$

Problem

actual
no end cap



$$R < d/2$$

$$d \ll L$$

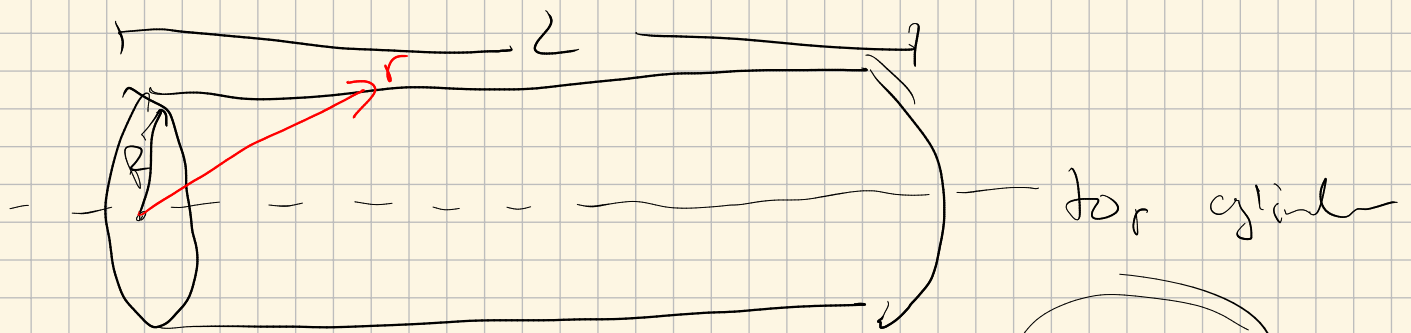
What is the capacitance
of this configuration?

1. Imagine putting $+Q$...
2. \vec{E} -field

3. ΔV between surfaces of the conductors

4. C ?

Electric field of one cylinder



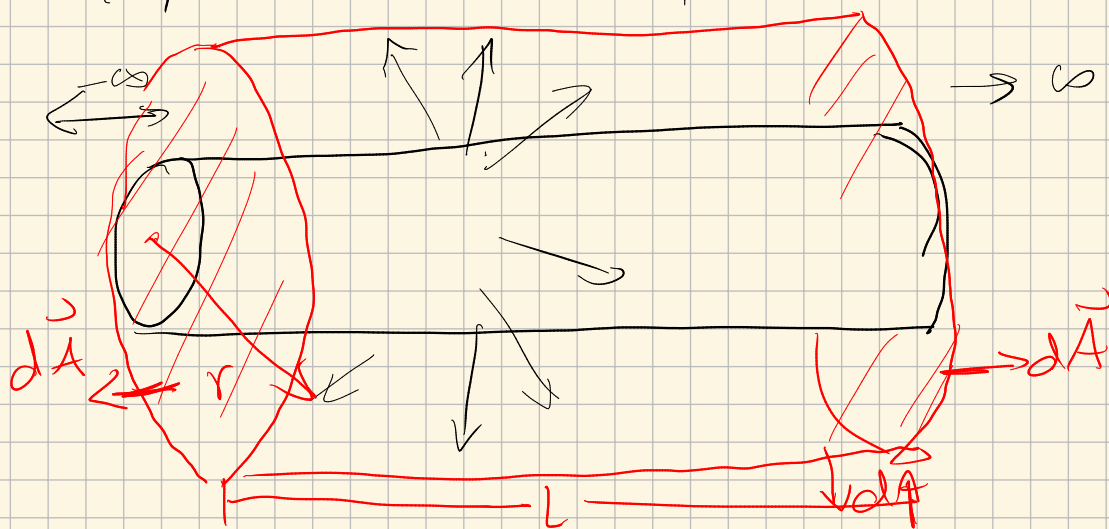
$$L \gg d$$

$$\vec{E} = \vec{E}_{\text{top}} + \vec{E}_{\text{bottom}}$$

approx. ∞ - long

(for || plate, $L \gg d$)

Gauss's Law for ∞ cylinder



$$\Phi = \int_{\text{closed surface}} \vec{E} \cdot d\vec{A} = \int_{\text{lateral surface}} \vec{E} \cdot d\vec{A} + \int_{\text{end caps}} \vec{E} \cdot d\vec{A}$$

$$\Phi = \int_{\text{lateral surface}} \vec{E} \cdot d\vec{A} \approx \int_{\text{lateral surface}} |E| dA$$

$$= |E| \int_{\text{lateral surface}} dA$$

area of lateral surface

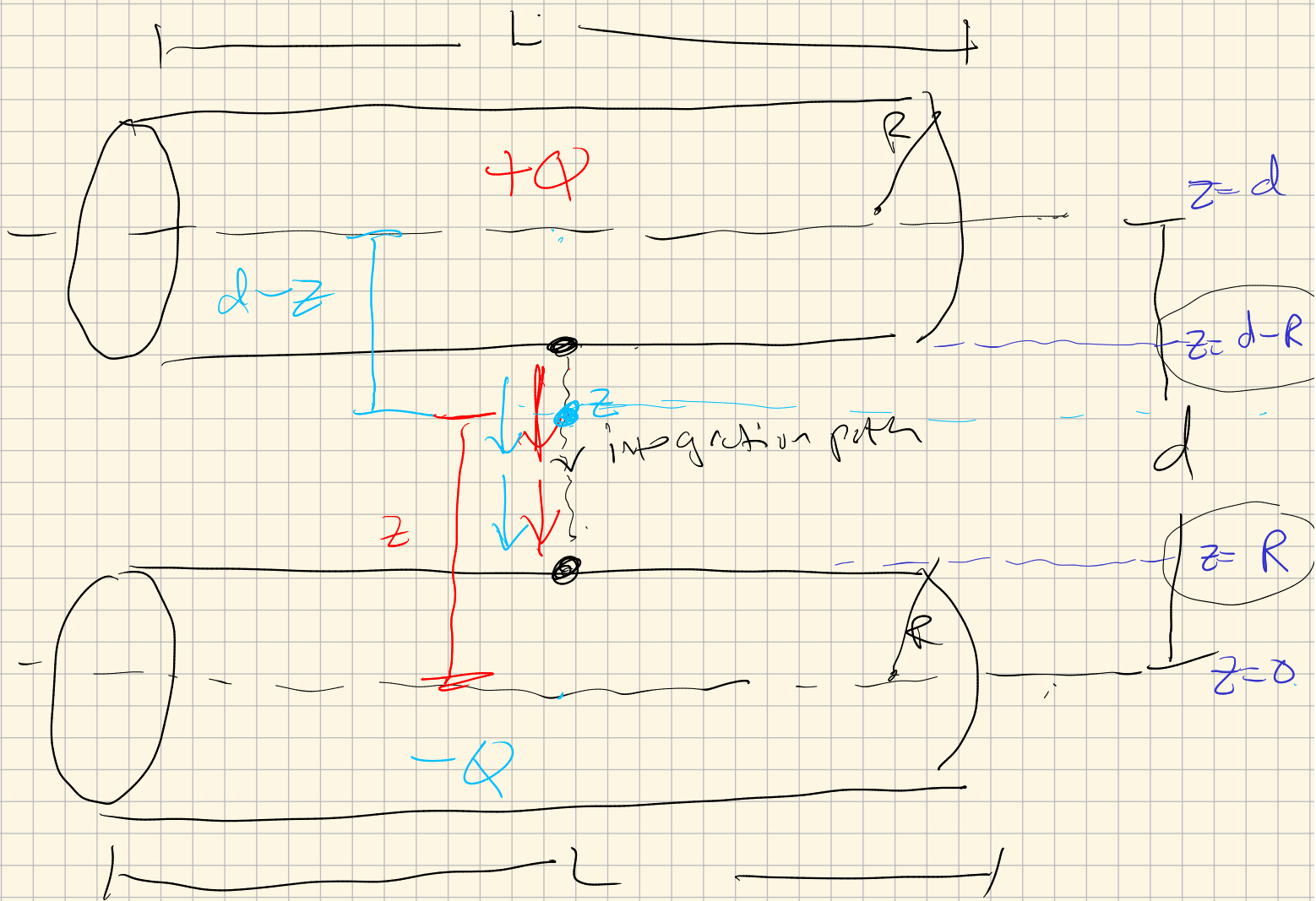
$$= |E| (2\pi r L)$$

$$\Phi = \frac{Q}{\epsilon_0}$$

$$|E| 2\pi r L = \frac{Q}{\epsilon_0}$$

$$|E| = \frac{Q}{\epsilon_0} \frac{1}{2\pi r L}$$

distance from the center of the cylinder



$$\vec{E}_{\text{total}} = \left(-\frac{Q}{2\pi(d-z)L\epsilon_0} - \frac{Q}{2\pi zL\epsilon_0} \right) \hat{k}$$

$$\Delta V = - \int_{z=d-R}^{z=R} \vec{E} \cdot d\vec{\ell} = - \int_{d-R}^R \left(-\frac{Q}{2\pi(d-z)L\epsilon_0} - \frac{Q}{2\pi zL\epsilon_0} \right) dz$$

$$= \frac{Q}{2\pi L\epsilon_0} \int_{d-R}^R \left(\frac{1}{d-z} + \frac{1}{z} \right) dz$$

$$= \frac{Q}{2\pi L \epsilon_0} \left[-\ln\left(\frac{d-R}{R}\right) + \ln\left(\frac{R}{d-R}\right) \right]$$

$$= \frac{Q}{2\pi L \epsilon_0} \ln\left(\frac{R^2}{(d-R)^2}\right)$$

$$= \frac{Q}{\pi L \epsilon_0} \ln\left(\frac{R}{d-R}\right)$$

$$C = \frac{Q}{|\Delta V|} = \frac{Q}{\frac{Q}{\pi L \epsilon_0} \ln\left(\frac{R}{d-R}\right)}$$

$$= \frac{\pi L \epsilon_0}{\ln\left(\frac{R}{d-R}\right)}$$

$R < d/2$
ex: $\frac{d/2}{d/2 - d/4}$

$$= \frac{\pi L \epsilon_0}{\ln\left[\frac{d-R}{R}\right]}$$