

Kramers-Kronig Relations

In the time domain, the linear polarization:

$$\bar{P}(t) = \epsilon_0 \int_{-\infty}^t \chi(t-\tau) \bar{E}(\tau) d\tau$$

Using the change of variables $\tau \rightarrow t-\tau$

$$\bar{P}(t) = \epsilon_0 \int_0^{\infty} \chi(\tau) \bar{E}(t-\tau) d\tau \quad (1)$$

This implies that the system is causal, i.e. $\chi(\tau) = 0$ for $\tau < 0$

Recall that the \bar{D} field is given by:

$$\bar{D}(t) = \epsilon_0 \bar{E}(t) + \bar{P}(t) = \epsilon_0 \bar{E}(t) + \epsilon_0 \int_0^{\infty} \chi(\tau) \bar{E}(t-\tau) d\tau \quad (2)$$

Taking the Fourier transform, we have:

$$\begin{aligned} \bar{D}(\omega) &= \epsilon_0 \bar{E}(\omega) + \epsilon_0 \int_{-\infty}^{\infty} dt e^{-i\omega t} \int_0^{\infty} \chi(\tau) \bar{E}(t-\tau) d\tau \quad (3) \\ &= \epsilon_0 \bar{E}(\omega) + \epsilon_0 \int_0^{\infty} \chi(\tau) \left(\int_{-\infty}^{\infty} dt \bar{E}(t-\tau) e^{-i\omega t} \right) d\tau \\ &= \epsilon_0 \bar{E}(\omega) + \epsilon_0 \int_0^{\infty} \chi(\tau) \left(\int_{-\infty}^{\infty} dt \bar{E}(t) e^{-i\omega(t+\tau)} \right) d\tau \\ &= \epsilon_0 \left[1 + \int_0^{\infty} \chi(\tau) e^{-i\omega\tau} d\tau \right] \bar{E}(\omega) \end{aligned}$$

Thus the complex permittivity is:

$$\epsilon(\omega) = \epsilon_0 \left[1 + \int_0^{\infty} \chi(\tau) e^{-i\omega\tau} d\tau \right] = \epsilon_0 [1 + \chi(\omega)] \quad (4)$$

Note that since $\chi(\tau)$ is real (time domain!), then the permittivity must satisfy

$$\epsilon(-\omega) = \epsilon^*(\omega)$$

which is evident from the form of (4). If we explicitly write $\epsilon(\omega)$ as a complex function:

$$\epsilon(\omega) = \epsilon'(\omega) + i\epsilon''(\omega)$$

We see that the real & imaginary parts must satisfy:

$$\epsilon'(\omega) = \epsilon'(-\omega) \quad \text{and} \quad \epsilon''(-\omega) = -\epsilon''(\omega)$$

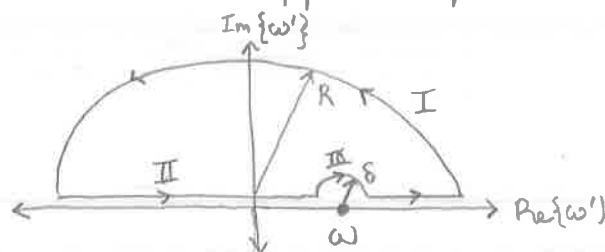
This means that the real part of the permittivity is an even function while the imaginary part is odd. Note that, in the limit that $\omega \rightarrow \infty$, we do not expect the material to respond (if only because charge has mass and thus cannot react infinitely quickly)

Thus $\epsilon(\omega)$ must have the property that $\epsilon(\omega \rightarrow \infty) \rightarrow \epsilon_0$ or $\chi(\omega \rightarrow \infty) \rightarrow 0$.

It turns out to be very informative to consider the integral:

$$I = \frac{1}{2\pi i} \oint_C \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega' = \frac{1}{2\pi i} \oint_C \frac{\chi(\omega')}{\omega' - \omega} d\omega'$$

where C is a contour in the upper complex ω' -plane:



Notice that $\chi(\tau)$ is finite and $\chi(\omega)$ is analytic. Since $\chi(\omega)$ is analytic, it follows that the whole integrand $\chi(\omega')/(\omega' - \omega)$ is analytic everywhere except at $\omega' = \omega$. Since the closed contour C does not enclose $\omega' = \omega$, the contour integral must be zero by the Cauchy Theorem

$$I = \frac{1}{2\pi i} \oint_C \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega' = 0$$

Next, as w/ a real space line integral, we can break this contour integral into integrals along the 3 connected open contours depicted above:

$$\int_{\text{I}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' + \int_{\text{II}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' + \int_{\text{III}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' = 0$$

Since $\mathcal{E}(\omega') - \mathcal{E}_0$ is finite, as we take $R \rightarrow \infty$, the integral along I must go to zero. This leaves us with two integrals: II is along the real ω' axis $-\infty < \omega' < \omega - \delta + \omega + \delta < \omega' < \infty$ and III is along a semi circle centered at $\omega' = \omega$ of radius δ

$$\int_{\text{II}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' + \int_{\text{III}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' = 0$$

Let us now consider the integral around contour III. If we make the substitution $\omega' - \omega = \delta e^{i\phi}$ then $d\omega' = i\delta e^{i\phi} d\phi$ and we have:

$$\int_{\text{III}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' = \int_{\pi}^0 d\phi \frac{\mathcal{E}(\omega + \delta e^{i\phi}) - \mathcal{E}_0}{\delta e^{i\phi}} i\delta e^{i\phi} = i \int_{\pi}^0 d\phi [\mathcal{E}(\omega + \delta e^{i\phi}) - \mathcal{E}_0]$$

↙ clockwise!

In the limit that $\delta \rightarrow 0$, the contour C encloses the whole upper ω' plane and we have:

$$\int_{\text{III}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega' = \lim_{\delta \rightarrow 0} -i \int_0^{\pi} d\phi [\mathcal{E}(\omega + \delta e^{i\phi}) - \mathcal{E}_0] = -i\pi [\mathcal{E}(\omega) - \mathcal{E}_0]$$

Thus it follows that

$$\mathcal{E}(\omega) - \mathcal{E}_0 = \frac{1}{i\pi} \int_{\text{II}} \frac{\mathcal{E}(\omega') - \mathcal{E}_0}{\omega' - \omega} d\omega'$$

The remaining integral over contour II is really just the principle part of the integral along the whole real ω' axis

That is:

$$\begin{aligned}\int_{\text{II}} \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega' &= \int_{-\infty}^{\omega - \delta} \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega' + \int_{\omega + \delta}^{\infty} \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega' \\ &= \text{PP} \int_{-\infty}^{\infty} \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega'\end{aligned}$$

Thus we are left with the relationship:

$$\epsilon(\omega) - \epsilon_0 = \frac{1}{i\pi} \text{PP} \int_{-\infty}^{\infty} \frac{\epsilon(\omega') - \epsilon_0}{\omega' - \omega} d\omega'$$

We can easily separate this into real and imaginary parts given $\epsilon(\omega) = \epsilon'(\omega) + i\epsilon''(\omega)$:

$$\begin{aligned}\epsilon'(\omega) - \epsilon_0 &= \frac{1}{\pi} \text{PP} \int_{-\infty}^{\infty} \frac{\epsilon''(\omega')}{\omega' - \omega} d\omega' \\ \epsilon''(\omega) &= -\frac{1}{\pi} \text{PP} \int_{-\infty}^{\infty} \frac{\epsilon'(\omega') - \epsilon_0}{\omega' - \omega} d\omega'\end{aligned}$$

These are the Kramers-Kronig relations.